

THE SEISMOLOGY PROGRAMME OF COROT.

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ABSTRACT

We introduce the main lines and specificities of the CoRoT Seismology Core Programme. The development and consolidation of this programme has been made in the framework of the CoRoT Seismology Working Group. With a few illustrative examples, we show how CoRoT data will help to address various problems associated with present open questions of stellar structure and evolution.

Key words: Stars: structure – pulsation – seismology – Space: photometry.

1. INTRODUCTION: THE SEISMOLOGY PROGRAMME AND THE COROT SWG

The main lines of the CoRoT seismology programme have been designed very early (Catala et al. 1995, Baglin et al. 1998) with strong specificities, in terms of requirements on the precision of frequency measurements, duration of the runs. Then, in order to develop this programme and optimize the scientific return, the CoRoT Seismology Working Group has been settled at the CoRoT kick-off meeting (1998, in Nice). The CoRoT SWG is intended to supply the required expertise and promote the investigation of relevant 'hard points'. This work is at the base of

the mission profile determination and of the targets selection. The SWG counts approximately 90 members in 20 institutes. The activity of the SWG has been regularly adapted to the needs associated to the evolution of the CoRoT project.

In Sect.2, we sketch out the scientific context for a project of stellar seismology like CoRoT. First (Sect.2.1), we try to give a flavour of the various scientific questions at stake. Then (Sect.2.2), we briefly come back on what has already been achieved so far and how this experience has logically lead the stellar international community to plan dedicated observations from space.

The CoRoT Seismology Core programme is described in Sect.3. For both Solar-like pulsators and "classical" pulsators (resp. Sect.3.1 and Sect.3.2), we present a few results obtained in the framework of the preparation of the CoRoT Seismology Programme. These results are selected to illustrate the preparatory work which has been made to investigate the various aspects of the Seismology Programme. The list is not pretended to be exhaustive and this is not our purpose either to develop these points in details. Some of them are developed further in this volume.

Finally, in Sect.3.3, we come back on a few aspects of the fields and targets selection process.

2. THE CONTEXT

2.1. Open questions of stellar physics and the seismology promiss

Stars are one of the main constituents of the Universe; they are also one of the major sources of information about it and thus an unavoidable subject of study. Nearly every field of astrophysics uses results of stellar structure and evolution theory, to estimate, for example, the age of globular clusters which give an essential piece of information on the age of the Universe, or to understand the origine of the chemical elements or the history of the Sun and of the solar system.

The main lines of stellar structure and evolution have been understood by confrontation of observables coming from the surface of stars and theoretical modelling calling to a wide panel of various fields of physics. Our understanding of stellar evolution and our capability to describe it precisely is thus suffering large uncertainties, due to the fact that a star is a complex object, involving a large number of physical processes still poorly understood.

Considering the upper part of the main sequence for instance, which is characterized by the existence of a convective core, one of the most debated open question is whether and how the central region mixed by

convection is extended by the so-called overshooting process. This point alone, by the change induced in the amount of hydrogen available for central nuclear reactions (see Fig. 1), is responsible for an uncertainty which can reach 30 to 50% in the age estimate of all stars with mass higher than $\sim 1.1M_{\odot}$ (for solar composition).

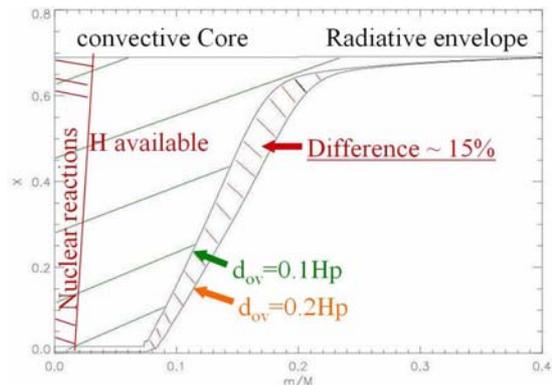


Figure 1. Hydrogen mass fraction profile (X) for a $1.8M_{\odot}$ stellar model at the end of Main Sequence evolution, for two illustrative values of the overshooting parameter d_{ov} . Hatched areas are proportional to the amount of hydrogen available for nuclear reactions on the Main Sequence, showing a difference of 15% between the two cases considered here.

For the lower part of the main sequence ($M \leq 1.4M_{\odot}$), characterized by an extended outer convective zone below the surface, one of the most prominent open questions deals with the efficiency of the heat transport in the upper convective regions. There, the density is so low that the heat transport efficiency depends severely on the description of the convective transport process. Since a fully consistent description of this process is still out of reach, the uncertainty in determining the temperature gradient becomes important and hampers severely our description of stars.

In spite of several tentative refinements, these two processes are widely considered at the moment in the modeling by simple one-parameter crude descriptions.

Beyond these two points, segregation of the different chemical species induced by gravitational and radiative forces are considered as a non negligible factor in several classes of objects, with a prominent manifestation in the surface anomalies observed in A stars. Impressive efforts have been made to implement these aspects (e.g. Michaud 2004), but they have been only confronted so far to surface classical observables and would benefit additional observational constraints.

Mass loss, meridional circulation and turbulence, their influence on angular momentum evolution,

their interaction with previously quoted diffusion of chemical elements constitute another active front for research in this field (e.g. Talon 2004, Vauclair and Théado 2003, Théado and Vauclair 2003), but here again, classical observables hardly can constrain such processes. On top of this, one can consider the effect of magnetic field and its potential interaction with the previous mechanisms (Alecian 2004, Mathis & Zahn 2005).

In the pre-Main Sequence phase, it is still an open question to know how the angular momentum evolves and what is its influence on the structure at the beginning of the Main Sequence phase.

At the other end of stellar evolution, the chemical composition profile of the strongly stratified structure of white dwarfs is holding the signature of e.g. $C^{12} - O^{16}$ poorly known reaction rate or mixing processes at work during the red giant phase.

In all these cases, stellar oscillations are expected to bring relevant additional constraints. Oscillations have been observed in stars representative of approximately all mass ranges and evolution stages, from the PMS stage, to the white dwarf cooling sequence, including Main Sequence, horizontal branch and red giant phases. New classes of stellar pulsators are still regularly discovered, and pulsation looks now more like the rule rather than the exception. This definitely suggests that seismology is a promising tool to improve our understanding of the various physical processes at work in stars.

Stars are generally seen pulsating on a more or less extended set of eigenmodes. These eigenmodes are standing waves established inside the stars by essentially two types of propagation waves: pressure waves, for which the restoring force is dominantly the pressure gradient, and gravity waves, for which the restoring force is essentially the buoyancy.

The associated eigenfunctions can generally be decomposed in a product of a function $f(r)$ describing the radial dependence and the angular dependence expressed in terms of spherical harmonics (see e.g. Unno et al. 89). Each eigenmode can thus be characterized by three integers: l and m , the degree and the azimuthal order of the spherical harmonics, and the radial order n , a count of the nodes of function $f(r)$.

The basic idea of stellar seismology is to use these new observables (frequencies, amplitudes, profiles, phase, ...) which are sensitive to the interior structure in a differential way, to help constraining the stellar structure and its evolution.



Figure 2. 'Et pourtant, elles pulsent!', Philippe Delache 1992.

2.2. From space and from the ground, state of the art of seismology

The case of our closest star is illustrative of the technique, though it has its own specificity. Observations of the Sun, from space and from the ground, have brought eigenfrequencies determined with precision up to a few nHz, thanks to very long and dense observation sequences (longer than 10 years with a 93% duty for SOHO, 20 years with Iris, 30 years with Bison, see e.g. Salabert et al. 2004; Garcia et al. 2004, 2005). The high spatial resolution allows to detect millions of eigenmodes. The results are impressive. The measurement of the rotation of the Sun with respect to radius and latitude down to the inner 20% central region (see e.g. Couvidat et al. 2003) has confirmed the existence of the tachocline, but has also assessed the solid rotation regime inside the radiative zone, which resists interpretation in terms of angular momentum transfer mechanism. Solar seismology has shown its ability to solve physical controversies in measuring precisely the central temperature, constraining to accept a neutrino mass (see e.g. Turck-Chièze et al. 2001, 2004).

Great efforts are developed to track gravity modes (e.g. Picard (see Thuillier et al. 2006), Golf-NG, see e.g. Turck-Chièze et al. 2005, 2006, see also Appourchaux et al. 2000) which would allow to probe the not yet resolved very inner part of the Sun.

Generalizing seismology techniques to all kind of

stars has been a leading idea for stellar astronomers. Though the accuracy and resolution obtained on the Sun look out of reach for distant stars, the large number of objects will allow a global vision of stellar evolution. Let us mention briefly the main development lines in this domain.

Solar-like oscillations are expected to exist in a large number of stars, in fact all F and early G stars, with amplitudes equal to or larger than in the Sun (see e.g. Houdek et al. 1999). These amplitudes remain intrinsically small (a few ppm in photometry and a few tens of cm/s in radial velocity) and the search for such oscillations in stars other than the Sun has been a long quest for the international community. After the first well-established detection about 7 years ago (Martic et al. 1999, Barban et al. 1999), oscillations have been found in a few more stars (see Bedding and Kjeldsen 2006 and references therein), thanks in particular to new spectrographs on large telescope like HARPS (see e.g. Mosser et al. 2004, Bouchy et al. 2005, Santos et al. 2004). These observations have provided tests for the theory of excitation of the modes in comparing predicted and observed amplitudes (Samadi et al. 2003 and ref therein). Interpretation of individual frequencies remains risky, due to the small signal to noise ratio and to the short duration of the observing runs. In the very favorable case of Alpha Cen A and B, for which oscillations have been found in both objects, several modelling studies have been made (Thevenin et al. 2002, Eggenberger et al. 2004, Miglio and Montalbán 2005), using seismic information and global parameters obtained from binarity characterization, and interferometry measurements. Miglio and Montalbán (2005), discuss the interest and the impact of the different observables on the study. In the actual state of the art however, it hardly leads to a firm conclusion in terms of physics, partly because, even in this case, most of the observable can still be questioned to some extent, but also because our experience in this domain still has to be built and secured by confrontation with better data for a larger set of objects.

For a large number of classical variable stars (δ Scuti, white dwarfs, sdB, Be, γ Dor,...) some modes show amplitudes larger than $\sim 10^{-3}$ in photometry (~ 100 m/s in radial velocity) and can be observed from the ground. Observations are planned regularly, often within coordinated multisite campaigns (STEPHI, WET, DSN,... see e.g. Li et al. 2004, Breger et al 2005, Vauclair et al. 2003, Mathias et al. 2004,...), which allow to reach these detection levels and to obtain satisfying time coverage and associated window function.

Interpretation in terms of modelling internal structure reaches different stages for different types of objects. For white dwarfs, various studies tend to determine the stellar parameters: total mass, mass of the envelope, rotation periode (see e.g. Pech et al. 2006) and the characteristic time scales on the cooling se-

quence. For δ Scuti and Be stars, fast rotation, very common in these classes of objects, has to be taken into account under several aspects of the modelling. This has driven theoretical developments, numerical implementation, and is now a crucial aspect of most of the present studies (Soufi et al. 1998, Lignières et al. 2003, Lignières et al. 2006, Reese et al. 2006).

For γ Dor objects, (which were not known 10 years ago), a coherent picture for the driving source of the oscillations has been reached only recently (Dupret et al. 2005). For all these objects however, it is fair to say that seismic interpretation has not yet reached the 'exploitation' level in terms of scientific return on physical processes at stake.

Space very early appeared as a predilection place for stellar seismology. In photometry, it is possible from space to track modes with amplitudes around 1ppm with very moderate apertures (30 cm for CoRoT and for objects with $m_v \sim 6$). Space also enables very high duty cycles and extended runs (up to 150d with a duty cycle higher than 90%, with CoRoT), giving access to characteristic time scales out of reach from the ground. The past two decades have seen an uninterrupted succession of proposals for national and international space projects dedicated to stellar seismology. The whole community is longing for such a unique point of view. The results from the Canadian experiment MOST (Matthews 2005) have raised an animated debate, illustrating the fact that the field is entering a new area. In a very close future (launch autumn 2006), CoRoT (Baglin et al. 2003) will constitute a major step in this domain, hopefully followed by other projects, like Siamois at Dôme C (see Bouchy et al. 2005, Mosser et al., this volume), Kepler, Plato (see Catala et al. this volume),...

3. COROT SEISMOLOGY CORE PROGRAMME - THE MAINLINES

The CoRoT Seismology Core Programme is addressing objects in a wide range of mass, between $1 M_{\odot}$ and $\sim 15 M_{\odot}$. It is highly focussed on main sequence evolution stage which represents 90% of the stellar lifetime. As commented herebefore, it is time now to make a qualitative step in the understanding and description of this stage of evolution. Several theoretical developments are proposed which need observational constraints. From early studies and as illustrated in examples given hereafter, it comes out that frequency measurements with precision of the order of $0.1 \mu\text{Hz}$, are sensitive to the detailed structure of stars in this evolution stage and susceptible to bring valuable discriminant tests for it. CoRoT seismology observation programme is thus intended to bring observational material of this quality for a significant set of objects. This is at the origin of one of the most specific characteristics of CoRoT: the possibility to dedicate long runs (up to 150 days) to the

same field.

Besides this guideline, an interest has been clearly affirmed for complementary shorter runs allowing an exploration of the pulsational behavior at the micro-mag level across the HR diagramme. The mission profile has thus been built around (at least) 5 long runs of 150 days each, completed with approximately the same number of shorter runs (~ 20 days).

3.1. CoRoT and Solar-like pulsators

The amplitudes observed in the Sun have been used as a dimensioning constraint to fix CoRoT specifications. This was initially justified by the theoretical prediction that amplitudes of solar-like pulsators increase with convective velocities (and thus with L/M), see e.g. Houdek et al (1999), and this has been confirmed since then by observations from the ground (see Samadi et al 2005). Solar like candidates for CoRoT have thus been selected among F stars on or near the Main Sequence. They are representative of masses between 1.1 and 1.5 M_{\odot} , where the structure is changing significantly from small or no convective core associated with an extended outer convective zone to an important convective core and a tiny envelope.

Hunting for a 0.1 μHz precision on frequency measurement. Several early prospective works have shown that eigenfrequencies, if determined with a precision of the order of 0.1 μHz were sensitive and discriminant in terms of physics options in the modelling of these objects. This is the case of Michel et al. (1995) who investigate the effect of slight variations of overshooting amount, mixing-length, metallicity, etc... This forward approach has been investigated further in the framework of 'Hare and Hounds' (H&H) exercises reproducing dimensioning factors of the expected CoRoT observations. Using simulated data of the H&H3 exercise for instance, Provost et al. (2002, see also Berthomieu et al. 2003, Provost et al. 2000), investigates the potential of classical frequency indexes $\delta\nu_{02} = \nu_{n,\ell=0} - \nu_{n-1,\ell=2}$ and $\delta\nu_{01} = 2\nu_{n,\ell=0} - (\nu_{n,\ell=1} + \nu_{n-1,\ell=1})$. As illustrated in Fig. 3, they concluded that $\delta\nu_{01}$ is a powerful indicator to discriminate main sequence and post main sequence evolution stages. More examples can be found in Appourchaux et al. (this volume) and Berthomieu et al. (this volume).

In order to explore what the final precision on frequencies determination could be, using the simulation tool developed by Baudin & Samadi (this volume), Michel et al. (2006) gave an illustration (see Fig. 4) of what kind of spectra and performances are expected for a solar-like target planned for a long run (150 days). Three cases are illustrated: a) the reference case: precision is estimated considering only photon noise and taking a 1 μHz generic value

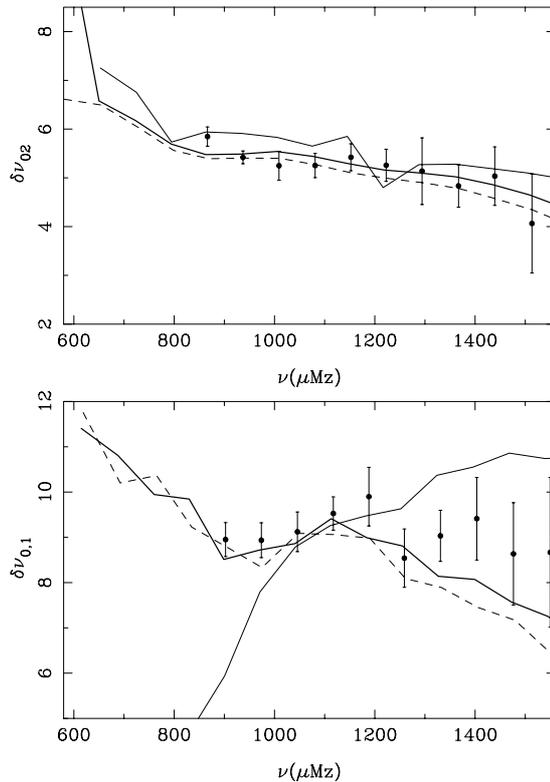


Figure 3. Small frequency spacings $\delta\nu_{02}$ and $\delta\nu_{01}$ as a function of the frequency, for two main sequence models M1 (heavy and dashed lines) and one post main sequence model (normal line) of HD 45067, compared to simulated data of HH3 exercise. Error bars result from the analysis of the simulated data.

for linewidth; b) as case a, but linewidth are from Houdek et al. (1999); c) same as b, but granulation noise contribution is considered in addition to photon noise. In order to illustrate the impact of the uncertainty on linewidth estimates, for case b and c, results are also shown for twice and half the values of the estimated linewidth. These results confirm the fact that the granulation noise as estimated here following Harvey (1985), might be a significant factor compared with photon noise. However, the expected precision on frequency determination remains of the order of a few 10^{-7} Hz, (below 0.7 μHz in the worst case considered here).

Tackling specific structural features One of the means to investigate the structure of the stellar interior from the oscillation frequencies without resorting to explicit modelling of the star is to utilise the oscillatory signal in the frequencies to determine the acoustic depth $\tau_d = \int_{r_d}^R dr/c$, of a sharp feature, lying at a radius of r_d , c being the sound speed and R the total radius of the star. This signal first commented by Gough (1990) can be amplified by using the second differences of the frequencies, which can then be fitted to a suitable oscillatory function to determine the acoustic depth, τ_d . Ballot et al. (2004)

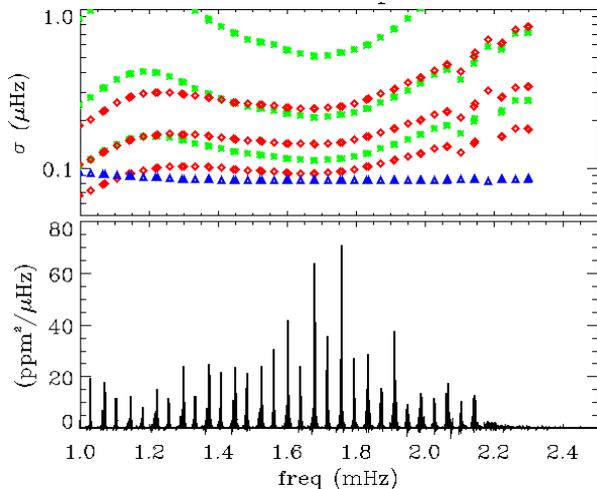


Figure 4. Lower panel: simulation of the pure seismic signal expected for HD49933. Upper pannel: Estimates of the 1- σ precision on the determination of eigenfrequencies, for case a, b and c (resp. triangles, rhombs, stars) as described in the text (from Michel et al. 2006).

have shown that the long runs of CoRoT would allow to extract the position of the bottom of the convective zone (BZC) of solar-like stars, within an accuracy of around 5% for the majority of solar-like targets.

Mazumdar (2005) applied this technique to the simulated CoRoT data for the primary target star HD49933 to correctly extract the acoustic depths of the base of the convective envelope and of the second helium ionisation zone of the input stellar model. Fig. 5 shows the functional fit to the simulated data with errors.

Such measurements with CoRoT data would constitute a strong constraint for better understanding convection in stars.

Rotation and inclinasion. Gizon & Solanki (2003) have studied the possibility of constraining both rotation rates (assumed to be rigid in first approximation) and stellar axis inclination from p-modes of stars spinning as slowly as two times the solar rate. Ballot et al. (2006) have investigated the value of using several modes simultaneously to increase the accuracy, especially at low angle (Fig. 6), and pointed out the strong correlation between the estimates of these two parameters.

Considering rotation profiles, and as stressed by the solar case, rotation inversion possibilities toward the center of the stars are very dependent on the detection of g or mixed modes. From the beginning, the idea thus has been to address this problem mostly with classical pulsators known to show this kind of

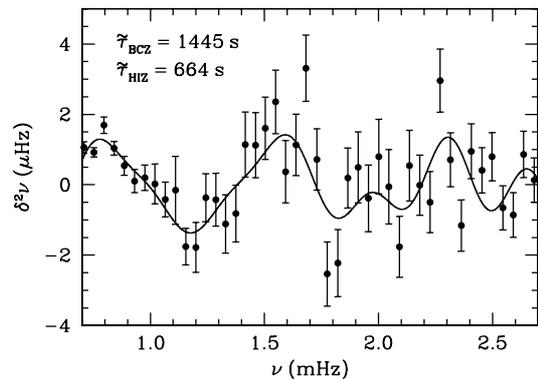


Figure 5. The oscillatory signal in the second differences of the (simulated) frequencies of HD49933 (shown as data points with respective errorbars) are fitted by a double-oscillatory expression (solid curve) to extract the acoustic depths of the base of the convective envelope (τ_{BCZ}) and the second helium ionisation zone (τ_{HIz}), (from Mazumdar 2005).

modes. However, Lochard et al. (2005) have shown (Fig. 7) that for appropriate solar-like targets observed with CoRoT, such modes being expected, it should be possible to have an estimate of the gradient in the rotation profile.

Amplitudes and convective transport. As shown by Samadi et al (2005), the study of amplitude distribution in this domain of the HR diagram can be used to constrain the properties of stellar turbulent convection. Indeed, the square of the mode amplitude, V^2 , is proportional to \mathcal{P} / η where \mathcal{P} is the rate at which energy is supplied by turbulent convection and η is the rate at which the mode is damped. Using several 3D simulations of stars, Samadi et al (2005) have found that the maximum of the excitation rate, \mathcal{P}_{max} , scales as $(L/M)^s$ where L and M are the luminosity and the mass of the star respectively and s is the slope of this power law. Furthermore the authors have found that the slope s is very sensitive to the way the convective eddies are time-correlated. Indeed, the slope s is equal to 3.2 when one models the eddy time-correlation according to a Gaussian function and to 2.6 when one models it according to a Lorentzian function. A comparison of their results using damping rates (η) from Houdek et al (1999), with available observations is strongly in favor of the Lorentzian description (see Fig. 8).

Diffusion processes. Several works have investigated the possibility to constrain chemical inhomogeneities induced by diffusion processes in stars (Théado et al. 2005, Castro and Vauclair 2006), confirming that with the accuracy expected for CoRoT, frequencies would constitute sensitive observables. In the case of solar-like stars hosting planets, it has been proposed to use seismology with CoRoT to distinguish between two possible scenari which are

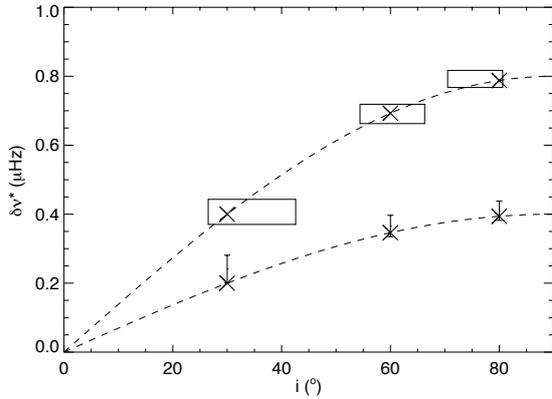


Figure 6. Biases and error bars for parameters i and $\delta\nu^* = \delta\nu \sin i$ obtained with a multi-mode fitting for 6 different simulated configurations ($\delta\nu_0 = \Omega/2\pi = 0.4, 0.8 \mu\text{Hz}$ & $i_0 = 30, 60$ and 80 degr). The crosses mark the expected values ($i_0, \delta\nu_0^*$). For $\Omega = 2 \Omega_\odot$ cases, the boxes indicate the mean results and their dispersions. For $\Omega = \Omega_\odot$ cases, only error bars on $\delta\nu^*$ are plotted because of the absence of good determinations of i . The two dashed lines are isorotations $\delta\nu = \delta\nu_0 = 0.4$ and $0.8 \mu\text{Hz}$ (from Ballot et al. 2006).

currently considered to explain metallicity excess in stars with planets (Bazot and Vauclair 2004, Bazot et al. 2005, see also Soriano et al. in this volume).

3.2. CoRoT and classical pulsators

Modal stability and convection description.

The delta Scuti, beta Cephei, gamma Doradus and PMS stars to be observed by COROT are auto-driven pulsators. The non-adiabatic modelling of stellar pulsation enables us to determine which modes are stable or overstable and to localize the

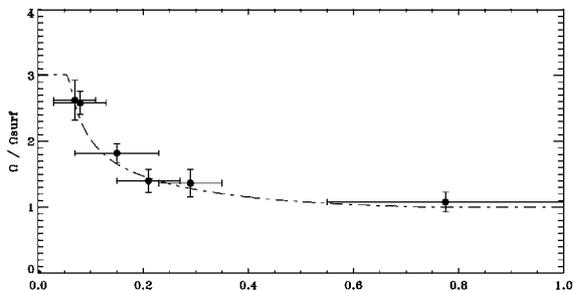


Figure 7. Dashed line: input profile for rotation rate (normalized to the surface here). Black dots and associated error bars: values obtained for inversion at different radii, considering representative noise and using rotational kernels computed with a trial stellar model showing large separation close to the ones of the input model. (from Lochard et al. 2005)

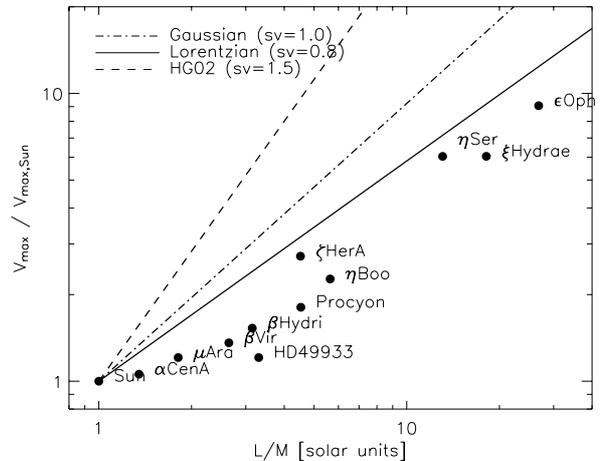


Figure 8. Maximum of the mode amplitudes (V_{\max}) relative to the observed solar ($V_{\max,\odot} = 33.1 \pm 0.9 \text{ cm s}^{-1}$) versus L/M (see text). Filled circles correspond to the few stars for which solar-like oscillations have been detected in Doppler velocity. The lines correspond to calculations obtained by Samadi et al. (2005) assuming Lorentzian eddy time-correlation functions (solid line) and Gaussian function (dot-dashed line). For comparison the dashed line shows the result by Houdek & Gough (2002), using a Gaussian function and classical Mixing Length Theory.

driving and damping regions inside the stars. In delta Scuti and gamma Dor stars, the description of convection and its time-dependent interaction with oscillations plays a major role in this driving; and in the beta Cephei stars, it is the metallic content and its location (where it accumulates due to transport mechanisms). Hence, the comparison with the observed excited modes enables us constraining these aspects. A specific analysis can be performed for each star, comparing the theoretical and observed range of overstable modes. But also, a general study can be made for each type of star, comparing the theoretical and observed instability strips in the HR diagram. As an illustration, we give in Fig. 9 the theoretical instability strips obtained for the gravity modes of gamma Doradus stars with the time-dependent convection (TDC) treatment of Grigahcene et al. (2005), for different values of the mixing-length parameter α (related to the size of the convective envelope) (Dupret et al. 2005). This illustrates how the description of convection can be constrained by a stability analysis. In this case, the best agreement is obtained for models with $\alpha = 2$.

Rotational profile inversion. As already mentioned, one of the assets of intermediate and high mass pulsators on the main sequence resides in the fact that they are expected to show mixed modes. As

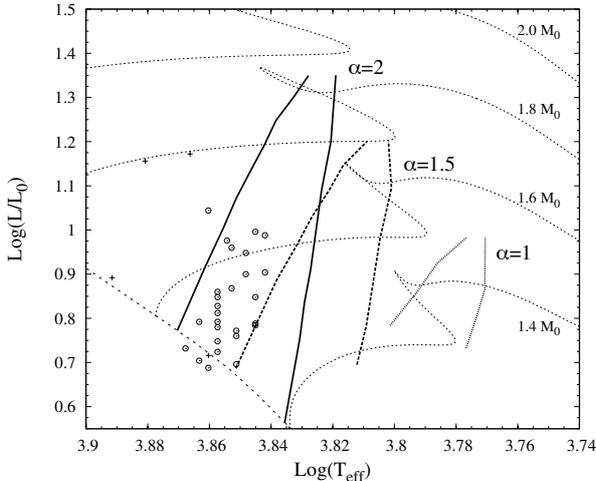


Figure 9. γ Dor theoretical IS for $\ell = 1$ modes, for three families of models with different values of α : 1, 1.5 and 2 obtained with TDC treatment (thick lines). The small circles correspond to observations of 27 bona fide γ Dor stars. $\alpha = 2$ theoretical IS best agree with observations.

shown by Goupil et al. (1996), for such objects, it is possible to build rotational kernels sensitive down to the very central regions and thus study the angular momentum transfer, one of the key aspect of this evolution stage for these objects.

Hot objects Beta cephei stars are main sequence stars with masses around $10 M_{\odot}$. Their structure is rather simple: a large convective core surrounded by a radiative envelope. Their metallicity is close to solar. Their oscillations are excited by the kappa mechanism, and they exhibit low degree and low order p and g modes. Their spectrum of oscillations is rather sparse and they are slow rotators, so that the multiplets due to rotational splitting can be identified. It has been shown from ground-based observations that asteroseismology of these stars can provide precise information on their global parameters, such as their mass, radius, age, metallicity, overshooting parameter, but also on their internal rotational law (Aerts et al. 2003). However, several problems remain unsolved. Indeed, some of the observed modes of oscillation cannot be excited using standard stellar models (Ausseloos 2004); they also present variable surface enhancements of nitrogen, which are hard to explain given that they are slow rotators (Morel et al. 2006). It is necessary to include non-standard physics to explain these observations.

A hare-and-hound exercise on beta cephei stars was done to prepare the COROT mission (Thoul et al 2003). The conclusion reached was that due to the simplicity of these stars it was possible to reconstruct the original star to a very high level of precision (Fig. 10). It was also concluded that in order to discriminate between different models, it was useful

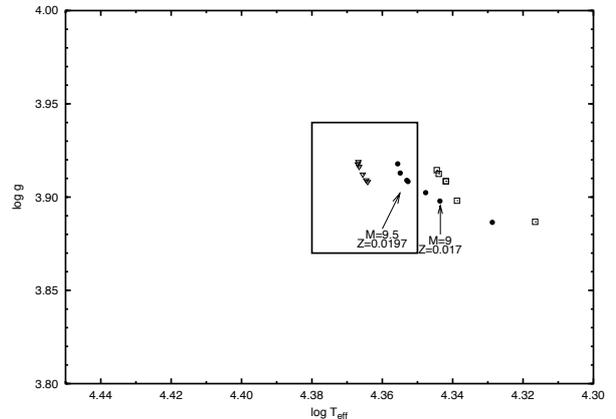


Figure 10. Positions of the stellar models which fit exactly two observed frequencies of oscillation of the β Cephei star HD 129929 in a $\log T_{\text{eff}} - \log g$ diagram. The squares, dots and triangles are obtained for models with $\alpha_{\text{ov}} = 0, 0.1, 0.2$, respectively. The observational error box from photometry is also given.

to observe more modes, including modes of degree $\ell > 2$, which are not observable from the ground. In addition, several multiplets have to be observed in order to probe different depths in the star if we want to get information about the internal rotation law.

In order to prepare the analysis of the data that COROT will provide on beta cephei stars, a database is being constructed, which contains stellar models and their oscillation frequencies. This database is described in details Thirion and Thoul (this volume).

g-modes pulsators. γ Doradus stars pulsate in the asymptotic g-mode regime. This makes it possible to obtain relevant physical information through the analytical expressions that the asymptotic theory provides. In this particular case, the eigenfrequency is given by (Tassoul 1980)

$$\sigma_{\text{asympt}} = \frac{\sqrt{\ell(\ell+1)}}{\pi(n+1/2)} \int_{r_a}^{r_b} \frac{N}{r} dr \quad (1)$$

where n is the radial order, ℓ the spherical order, r_a and r_b are the lower and upper limits of the radiative envelope of this stars and N is the Brunt-Väisälä frequency.

Therefore, as suggested by Moya et al. (2006), the ratio between two frequencies in the asymptotic regime depends only on the radial and spherical orders, taking the form

$$\frac{\sigma_1}{\sigma_2} = \frac{n_2 + 1/2}{n_1 + 1/2} \quad (2)$$

This allows us to estimate the radial order of observed γ Doradus frequencies. With at least three observed frequencies we can infer, through this procedure, some possible values of the radial order of each frequency. Once an estimate is fixed, the asymptotic expression provides a value for I the Brunt-Väisälä integral in (1). This gives us a new observable to be fitted by models.

If we display, for a given observed star, the estimated integrals as a function of the effective temperature, a figure giving a new constraint for the modeling of these stars is obtained (see Fig. 11 for an example in the particular case of 9 Aurigae). This technique has been successfully used for different γ Doradus

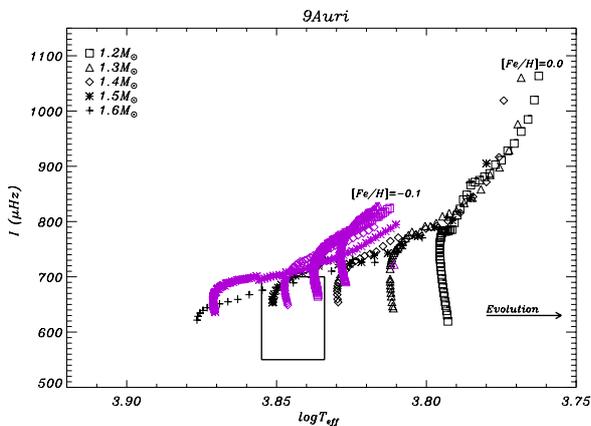


Figure 11. I - T_{eff} diagramme (cf text). The box represent the estimated value of the integral from observations, and the dots are the model predictions for different masses, metallicities, overshootings and evolutionary stages. From Moya et al. 2006.

Time/frequency analysis. For several class of classical pulsators, the question of the variability of the mode amplitudes is still an open question lacking seriously well suited observations. CoRoT long runs will offer a unique opportunity to apply time/frequency analysis and address this question. In order to get an estimate of the precision that one can expect for amplitude variations determination, F. Baudin made the simulation illustrated in Fig. 12: The application of time/frequency analysis to a 150 day long simulation of a sinusoidal oscillation of constant amplitude (400 ppm, i.e. representative of a lowest limit of what is seen from the ground) at the frequency $\nu = 100 \mu\text{Hz}$. In this simulation, a $m_V=8$ star is considered. The apparent power variations are due to the presence of noise (including activity and granulation). Their standard deviation corresponds to a 3% variation in power, for a time resolution of 5 days. Of course, this precision on the power variations will vary with the chosen time resolution of the analysis.

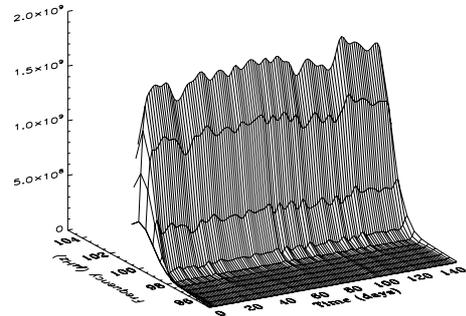


Figure 12. Time/frequency analysis of a 150 day simulation of a sinusoidal oscillation of constant amplitude (400 ppm) at the frequency $\nu = 100 \mu\text{Hz}$. Photon noise, granulation and activity signals have been computed using the simuLC tool described in Baudin et al. (this volume)

3.3. CoRoT sample:

Among the criteria used to select targets, observational performances have of course been considered. In this respect, a rapid estimate shows that for classical pulsators, in the hypothesis of stable frequencies and amplitudes over 5 months, a gain by a factor 500 at least can be made down to $m_V=9$ in terms of S/N, compared with what is currently achieved from the ground (see e.g. Michel et al. 2006). This lets room to choose these targets according to other criteria, like evolution stage, rotation rate, binarity,...

In the case of solar-like pulsators, the situation is less 'comfortable' and apparent magnitude always appears as a high priority parameter for selection. This is why these objects appeared very early and with a strong priority in the process of fields selection. First, the so-called Principal candidate stars have been selected with the conservative criterion based on observed solar amplitudes. Then, three criteria have been defined (Samadi et al. 2004) to evaluate the interest of secondary candidates to be selected in the field around. These criteria are intended to determine for which objects, a 'significant' number of modes can be expected to be measured with a given minimal precision. The amplitudes are estimated following Samadi et al. (2005) for different values of the linewidths, as commented hereafter. The noise level is obtained considering photon noise for CoRoT. A 150d duration of the run is assumed. Then, following Libbrecht (1992), an estimate of the frequency precision that could be obtained for a peak with half the maximum expected amplitude is derived.

Criterion 1 is our standard. It is referring to objects for which, assuming a generic $2 \mu\text{Hz}$ linewidth

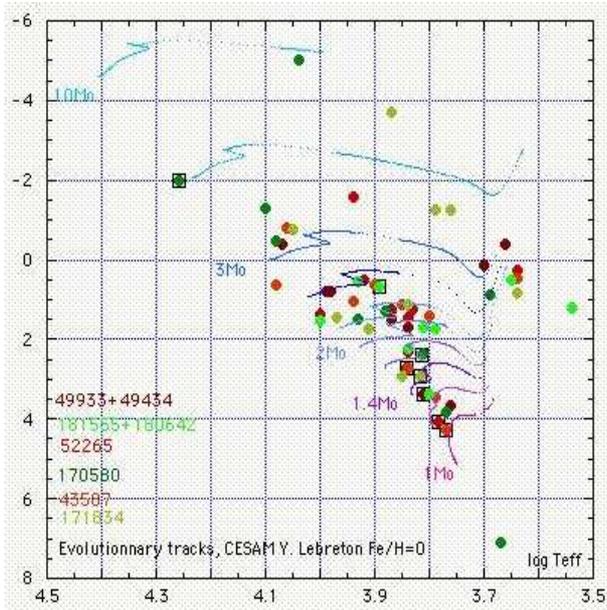


Figure 13. HR diagramme built with an illustrative selection of candidates potentially observed during 150 days with CoRoT. From Michel et Baglin (2005).

(~ 1.8 days lifetime), modes with amplitudes higher than half the expected maximum (i.e. a significant amount of peaks) would have frequency measured with precision better than $0.25 \mu\text{Hz}$ (associated with a signal to noise ratio $S/N=4$, as defined by Libbrecht (1992)). The application of this test is illustrated on Fig. 14,

Criterion 2 is an extrapolation of criterion 1 in the pessimistic eventuality of a $5 \mu\text{Hz}$ linewidth (~ 0.7 days lifetime), For objects satisfying this test, modes with amplitudes higher than half the expected maximum (i.e. a significant amount of peaks) would still have frequency measured with precision better than $0.4 \mu\text{Hz}$ (associated with a signal to noise ratio $S/N=4$, as defined by Libbrecht (1992)).

Criterion 0, the lowest one, also assumes a $2 \mu\text{Hz}$ linewidth, but corresponds to a $S/N=1$ value, which would allow the detection of a significant number of oscillation peaks, but is not expected to bring very precise frequency values. Roughly all stars brighter than $m_V=7.5-8$ satisfy this criterion.

Michel and Baglin (2005), with a preliminary selection of potential targets gave a flavor of how the CoRoT sample of stars observed during long runs could distribute in an HR diagramme. As shown in Fig. 13, it is possible to obtain a reasonable scan of the domain of interest in the HR diagramme. An updated picture of the target selection state of the art can be found in Michel et al. (this volume).

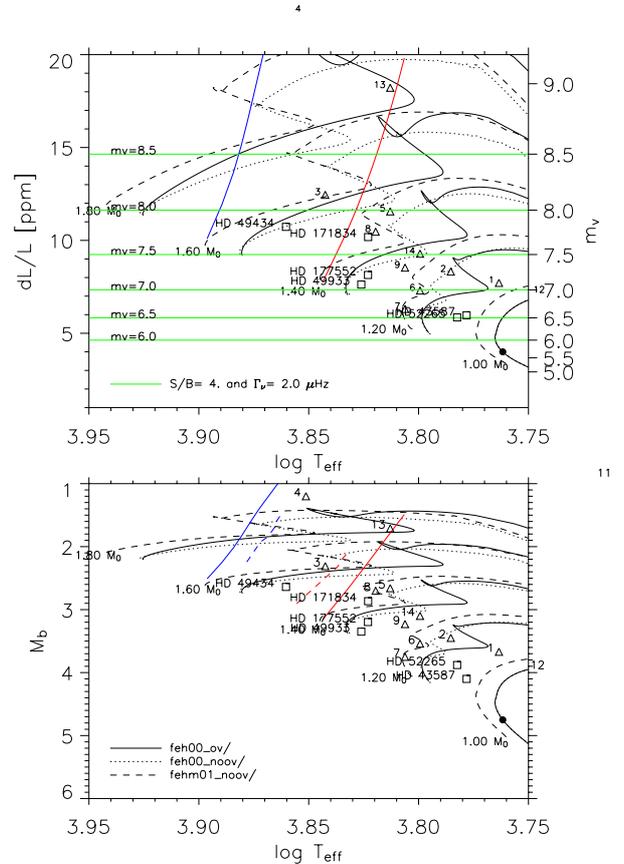


Figure 14. criterion 1. Candidate targets are attributed a mass and a luminosity by considering their location in the HR diagramme (lower panel) compared with evolution tracks. In the upper panel, the same evolution tracks are put in a diagramme with maximum amplitude estimates in ordinates (left axis). On the right axis are given values of observational magnitude m_V . Each of them (e.g. 7.0) separates the diagramme in two parts. If a candidate target is found in the upper part (above the corresponding green line), and if its observational magnitude m_V is lower than the limit value associated with this line (here 7.0), then the target satisfy the criterion. For instance candidate 1 would satisfy the present criterion if it is brighter than $m_V=7.0$.

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